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Held and Suarez (1994)再訪

今日の発表 とりあえずのShow and Tell (こんなのやりました)

- ・イントロダクション
 - ・Held and Suarez 1994の簡単な紹介
 - 私の動機
- 実験結果とHS94の図との比較など
- ・遭遇したり、思いついたりした、いくつかの課題
- ・ まとめ



イントロダクション: Fleid and Suarez 1994 の簡単な紹介」と 「私の動機」



Held and Suarez 1994 が目指すところ 1 ・「力学コア」に絞った「気候」モデル相互比較のためのテスト(基準/ベンチマ

- AMIP (フルAGCMsの気候モデル相互比較:Gates 1992) はモデルが複雑 過ぎる
 - 各モジュールの相互比較も必要
 - ICRCCM (放射コード) :Ellingson et al. 1991; Fouquart et al. 1991
 - ・陸面モデル: Henderson-Sellers et al. 1992)
- ・Williamson et al. (1992)は2D浅水波の「気象」(その後のWilliamson関連 テストなども「気象」(のはず))

ーク)の提案(「力学」と「気候 = 気象の統計状態」)

Held and Suarez 1994 目指すところ 2

- ・モデルのモジュール化の促進も
 - 取っ替え引っ替え
 - 物理過程についてはKalnay et al. (1989)が議論
- ・モデルの公開も

問題点

- ・正解が分かっていない
- ・長期変動がもたらすサンプリング・エラー





何を 今さら HS94? ほかにもいろいろ応用されてきたが、原点に近いところをなぜ?

- cuss their sensitivity to the choice of **dissipation**.
 - 形跡がない…

 - ・スペクトル・モデル同士の比較だと、極めてつまらない?
 - まあ、それでも、とりあえず、ちょっと何かやってみるかと思った

 In a separate study, we will be looking at this proposed calculation in more depth. In particular, we will study the behavior of these two models as a function of horizontal and vertical resolution and dis-

本人たちもやったかもしれないが(というかやったはずだが)出版された

・おそらく、他の研究者がそれなりにやってしまった(レビュー必要)?

何を 今さら HS94? e.g. Wan et al. 2008 (ECHAM5, T31-159, L16-81) • Results show that the simulated climate state in the Held–Suarez test is sensitive to spatial resolution. Increase of horizontal resolution leads to slight weakening and poleward shift of the westerly jets. Significant warming is detected in high latitudes, especially near the polar tropopause, while the tropical tropopause becomes cooler. The baroclinic wave activity intensifies considerably with increased horizontal resolution. Higher vertical resolution also leads to stronger eddy variances and cooling near the tropical tropopause, but equatorward shift of the westerly jets. The solutions show an indication of convergence at **T85L31** resolution according to all the three statistical tests applied. Differences between integrations with various time steps are judged to be

・とりあえず、たぶん、これの後追いに近いことから…

- within the noise level induced by the inherent low-frequency variability.





Explicit global simulation of the mesoscale spectrum of atmospheric motions

Yoshiyuki O. Takahashi,¹ Kevin Hamilton,² and Wataru Ohfuchi³ these experiments is shown in the upper right panel. We find that to obtain the convergent spectra we need to scale the diffusivity approximately as a power law of the model truncation (the regression fit shown in Figure 3 is $1.2 \times 10^{21} n_t^{-3.22} \text{ m}^4 \text{s}^{-1}$, where n_t is the truncation wavenumber).

Figure 3. As in Figure 2 but for AFES run with different numerical resolution. Results are shown for the 24 level version truncated at T79, T159, T319 and T639, as well as the T639L48 version. At each horizontal resolution a diffusion coefficient has been determined by trial and error to produce the fairly convergent behavior at the high wavenumber end of the spectrum. The black symbols in the inset show the diffusion coefficient as a function of truncation obtained this way. The red dots show results from a similar analysis of a version of the NCAR atmospheric model obtained by *Boville* [1991]. The lines in the inset are linear regressions.



実験とHS94の図との比較など

実験の強制・設定など Held and Suarez (1994)

力学コアのみ

乾燥大気、地形なし

地球的パラメータ

DCPAM5-20180304-2

水平粘性



・最高波数のe-folding time: 0.1 [day]

$$\frac{\partial v}{\partial t} = \dots - k_v(\sigma) v$$

$$\frac{\partial T}{\partial t} = \dots - k_T(\phi, \sigma) \left[T - T_{eq}(\phi, p) \right]$$

$$T_{eq} = \max \left\{ 200 \text{K}, \left[315 \text{K} - (\Delta T)_y \sin^2 \phi - (\Delta \theta)_z \log \left(\frac{p}{p_0} \right) \cos^2 \phi \right] \left(\frac{p}{p_0} \right)^{\kappa} \right\}$$

$$k_T = k_a + (k_s - k_a) \max \left(0, \frac{\sigma - \sigma_b}{1 - \sigma_b} \right) \cos^4 \phi$$

$$k_v = k_f \max \left(0, \frac{\sigma - \sigma_b}{1 - \sigma_b} \right)$$

$$\sigma_b = 0.7 \qquad k_f = 1 \text{ day}^{-1},$$

$$k_a = \frac{1}{40} \text{ day}^{-1} \qquad k_s = \frac{1}{4} \text{ day}^{-1}$$

$$(\Delta T)_v = 60 \text{K} \qquad (\Delta \theta)_z = 10 \text{K}$$

$$p_0 = 1000 \text{ mb} \qquad \kappa = \frac{R}{c_p} = \frac{2}{7} \qquad c_p = 1004 \text{ J kg}^{-1} \text{ K}^{-1}$$

$$\Omega = 7.292 \times 10^{-5} \text{ s}^{-1} \qquad g = 9.8 \text{ m s}^{-2} \qquad a_e = 6.371 \times 10^6 \text{ m}.$$





- ・解像度(σ等間隔)[Δt (min.)]
 - L20
 - T10 [60], T21 [60], T42 [30], T85 [15], T170 [6]
 - T42
 - L320 [4]
- •初期条件:等温(250 K + ノイズ)静止
- 6時間ごと出力
- ・ 各1200日積分, 解析は最後の1000日

• L5 [30], L10 [30], L20 [30], L40 [30], L80 [20], L160 [5],





強制温度 Temperature and Potential Temperature: Figure 1 a b



- 330 - 290 - 270

T42L20 vs HS94 U: Figure 2



Held and Suarez (1994)

- 36

- 24

- 12

- 0

- -12

- -24

-36



FIG. 2. The zonal-mean zonal wind produced by the T63 spectral model and G72 gridpoint model. Both are 1000-day means. Since the forcing is symmetric about the equator, differences between the hemispheres are indicative of sampling errors.

何も考えずに水平解像度を変える:





何も考えずに水平解像度を変える:



何も考えずに鉛直解像度を変える:







T42L20 vs HS94 (G72) Temperature and Potential Temperature: Figure 1 c d





何も考えずに水平解像度を変える:







na

sigr



0.0

0.2

0.4

0.6

na

sig

- 315

- 300

- 285

- 270

- 255

-240

- 225

-210

- 195

180



T85L20



Latitude [deg]

T21L20

Zonal Mean Temp [K] : hs94: T170L20: : : time = 200.0000 - 1200.0000 [day]



T42L20

T170L20



何も考えずに鉛直解像度を変える:





T42L5













T42L20 vs HS94 T* Variance: Figure 3



50

40

- 30

- 20

- 10



Held and Suarez (1994) Fig. 3. As in Fig. 2 but for the eddy variance of the temperature.

何も考えずに水平解像度を変える: T* Variance





- 50

40

- 30

- 20

- 10



Variance (Temp [K]) ** 2: hs94: T42L20: : : time = 200.0000 - 1200.0000 [day] 0.0 0.2 0.4 0.6 0.8 1.0 -90 -60 -30 0 30 60 Latitude [deg]

T42L20







何も考えずに鉛直解像度を変える: T* Variance







T42L80



T42L20







T42L20 vs HS94 U*^2 Spectra: Figure 4



HS94と比較して値が1.5~2倍くらい? フーリエ係数の定義等要チェック

Held and Suarez (1994)





FIG. 4. As in Fig. 2 but for the vertically averaged zonal spectra of the eddy variance of zonal wind.

何も考えずに水平解像度を変える: U*^2 Spectra



Zonal Power Spectrum (U [m/s]) ** 2 : hs94: T85L20: : : time = 200.0000 - 1200.0000 [day]





T85L20

Zonal Power Spectrum (U [m/s]) ** 2 : hs94: T170L20: : : time = 200.0000 - 1200.0000 [day]

T170L20



何も考えずに鉛直解像度を変える: U*^2 Spectra



T42L40

T42L80

何も考えずに)水平解像度を変える: V*Temp*





T10L20





-8

-16

-24

- 24

16



T85L20

0.2 0.4 gm 0.8 1.0 -90 -60 -30 30 60 Latitude [deg]

T42L20

T21L20

Zonal Mean V* x Temp* [K m/s] : hs94: T170L20: : : time = 200.0000 - 1200.0000 [day]

T170L20





何も考えずに鉛直解像度を変える: V*Temp*

- 24





T42L5





T42L80

T42L10

T42L20

T42L160

何も考えずに)水平解像度を変える: **V*U***

72

48





T10L20





T170L20

何も考えずに鉛直解像度を変える: **V*U***



T42L5

-30





T42L80



T42L10



T42L20







遭遇したり,思いついたりした,いくつかの課題

いくつかの課題

- ・地表面付近の温度構造問題
 - ・地表面付近が冷えすぎる
 - ガタガタ問題
- 長期変動
- ・水平粘性 / エネルギー・スペクトル(再訪)



地表面付近に逆転層? 107本





長期変動 経度平均Uの時間変動

Time [day]

Spatial structure of ultra-low-frequency variability of the flow in a simple atmospheric circulation model

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(Received 27 March 1991; revised 9 June 1992)

SUMMARY

The internally generated variability of the global circulation on time-scales in excess of one year is investigated using a multilevel baroclinic primitive-equation model with moderate resolution of dynamical processes. The model includes idealized heating and friction. Maximum variability is found on time-scales of around a decade. Empirical orthogonal-function analysis reveals that the principal component of this variability is an alternate splitting and coalescence of the subtropical and mid-latitude jets. Associated changes in the strength and distribution of eddy activity are found. The character of the ultra-low-frequency variability is completely changed for very low resolutions, but seems relatively insensitive to details of model formulation at higher resolution. The possibilities of deriving a low-order model which includes ultra-low-frequency variability are considered.

551.513.1

Figure 3. Ultra-low-frequency variability of the atmospheric angular velocity in Run 1. (a) Time series of the angular velocity from year 10 to year 58 of the run. (b) Amplitudes of the Fourier coefficients for the time series from year 4 to year 100.

- ・解像度を変えたとき、粘性係数をどう変えれば良いのか?
- ・超粘性の次数への依存性

散逸(水平粘性など)とエネルギースペクトル

Held and Suarez (1994)再訪:まとめ ・とりあえずDCPAMを走らせることができた(京大数理研竹広さんに感謝)

- ・テスト計算でHS94実験を解像度を変えて行った(網羅的ではない)
- スペクトル + 水平粘性などの問題が挙げられる
- 用), TEM解析]
- 態ではないが)

 統計量は解像度に依存し、解像度を上げると収束しているように見えなくも ないが、散逸の調整をきちんと考えていないので、更なる検討が必要である

・そのほか,今後の課題として,地表面付近の温度構造,長期変動,エネルギー

・Pythonで解析と描画ができるように整備中 [特に,内挿(σ→p座標変換を含 む),スペクトル解析(松嶋さんのISPACKのPythonインタフェースを活

・Williamson et al. (1998)の拡張(WOB98)は行った(他人に見せられる状

