





i) Dry magmas ascend up a conduit very much like long gravity currents, under an interplay of fluid mechanics, thermodynamics and elasticity.a) Fluid mechanics and thermodynamics

8.3











8.8
Rock elasticity
$$Vdp = \beta_r dV$$
 (β_r rock bulk modulus ~ 10¹⁰ Pa) (*)
 $V = \frac{1}{\beta_r} + \frac{1}{\rho} \frac{\partial \rho}{\partial p} \frac{dp}{dt} = \frac{Q}{\rho} - \frac{V}{\rho} \frac{\partial \rho}{\partial T} \frac{dT}{dt}$ (4)
 $(V/\beta_{eff}) \frac{dp}{dt}$ effective thermal (**)
Solubility $n = N - sp^{1/2}(1 - x)$ (5)
n mass fraction of exsolved volatiles; *s* solubility constant ~ 3 x 10⁻⁶ Pa^{-1/2}
Density relationship $\rho = \frac{nPT}{p} + (1 - n) \frac{x}{\sigma_c} + \frac{1 - x}{\sigma_m} \int_{m}^{-1} (6)$
(gas) (regstation (meth))







Equations for steady homogen	eous flow in pipe	8.12
mass conservation	$\rho uA = Q$	(1)
momentum conservation	$\rho u \frac{du}{dz} = -\frac{dp}{dz} - \rho g - f_{\text{friction}}$	(2)
density	$1/\rho = (1 - n)/\rho_s + nRT/p$ solid + gas	(3)
volatile content	$n = n_0 - sp^{1/2}$ Henry's law constant	(4)
friction	$f \sim \frac{1}{12} \mu u / r^2$ (viscous liquid)	(5a)
	~ $0.001 \rho u^2/r ~(\sim 0)$ (turbulent gas)	(5b)
void fraction	$\phi = 1 + \frac{(1-n)p}{nRT\rho_s}^{-1}$	(6)
	fragmentation at about = 75%	

8.13
(1), (2), and (3)
$$\frac{dp}{dz} 1 - \frac{u^2}{a^2} = -\rho g - f$$
(7)
where sound speed $a^2 = dp/d\rho = a^2(p)$ (8)
with $a \ 0.95 (n_0 RT)^{1/2}$ (9)
Integrate equations from
 $p = p_0$ ($z = 0$) (in chamber) (10)
 $p_e = p_a \ OR \ u = a \ (p_e > p_a) (z = H)$ (at surface) (11)
exit pressure





8.16 mass conservation $\frac{\partial}{\partial t} \rho + \frac{\partial}{\partial z} \rho w = 0$ w(z,t): vertical velocity momentum conservation $\frac{\partial p}{\partial z} = -\rho g - \frac{8\mu w}{r_E^2}$ crystal growth $\frac{\partial x}{\partial t} + w \frac{\partial x}{\partial z} = 4\pi \ \phi r^2 = (36\pi\phi)^{1/3} \ x^{2/3}$ number density of crystals $\phi(z, t)$; crystal radius r(z, t); constant linear crystal growth rate boundary conditions z=0 $\frac{dP}{dt} = \frac{\beta}{V} (Q_1 - Q_0); \ x = x_0$ z=H $p=p_{atm}$













Lecture 8. Explosive Volcanism Bruce, P.M. and Huppert, H.E., 1989 Thermal control of basaltic fissure eruptions, *Nature*, 342, 665-667. Huppert, H.E., Sparks, R.S.J. and Turner, J.S., 1982 Effects of volatiles on mixing in calc-alkaline magma systems, *Nature*, 297, 554-557. Huppert, H.E., and Woods, A.W., 2002 The role of volatiles in magma chamber dynamics, *Nature*, 420, 493-495. Melnik, O. and Sparks, R.S.J., 1999 Non-linear dynamics of lava dome extrusion, *Nature*, 402, 37-41. Morton, B.R., Taylor, G.I. and Turner, J.S., 1956 Turbulent gravitational convection from maintained and instantaneous sources, *Proc. Roy. Soc.* A 234, 1-23. Turner, J.S., Huppert, H.E. and Sparks, R.S.J., 1983 An experimental investigation of volatile exsolution in evolving magma chambers, *J. Volcanol. Geotherm. Res.*, 16, 263-277. Woods, A.W. 1995 The dynamics of explosive volcanic eruptions. *Rev. Geophys.*, 33, 495-530. Woods, A.W. and Huppert, H.E., 2003 On magma chamber evolution during slow effusive eruptions, *J. Geophys. Res.*, 108(B8), art. no. 2403.

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